

# Strangeness Production and $YN$ Interaction

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## Possibilities:

- ELASTIC  $\Lambda N$  SCATTERING

PROBLEM: SHORT LIFE TIME OF  $\Lambda \Rightarrow$

SCARCE DATA SET EXISTS

(G. Alexander et al., Phys. Rev. **173** (1968) 1452;  
B. Sechi-Zorn, et al., Phys. Rev. **175**, 1735 (1968))

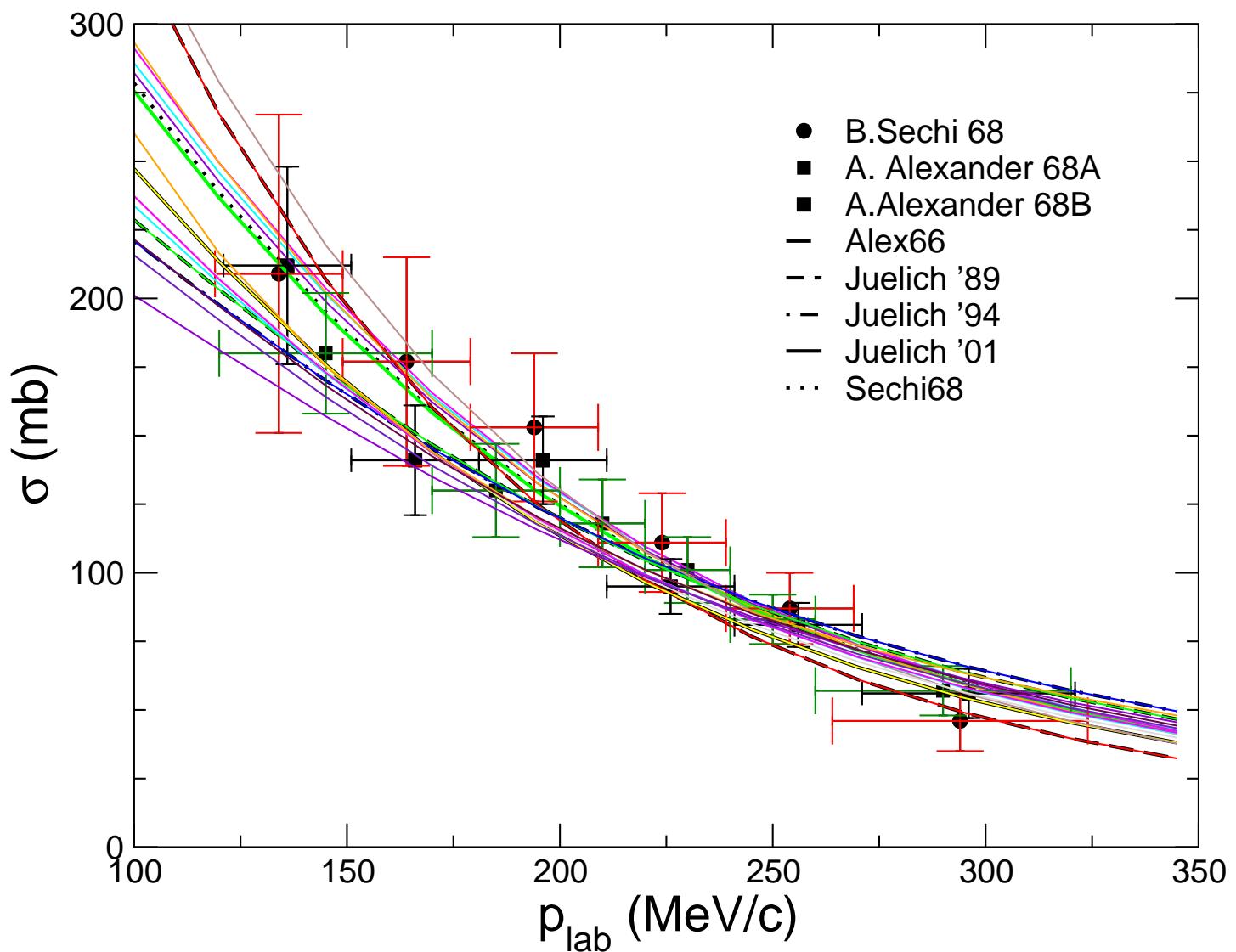
ONLY A FEW HUNDRED EVENTS

$\Rightarrow$  SCATTERING PARAMETERS  
ONLY BADLY DETERMINED !

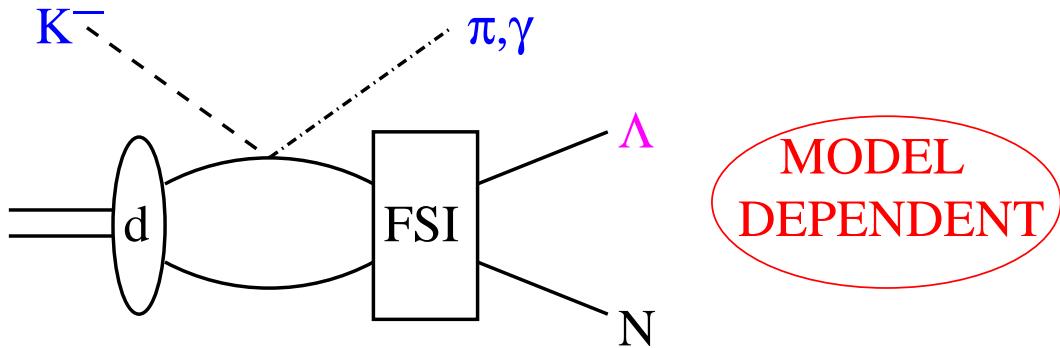
$$a(S = 0) = -0.7 \div -2.6 \text{ fm}$$
$$a(S = 1) = -1.7 \div -2.15 \text{ fm}$$

(V.G.J. Stoks, T.A. Rijken Phys. Rev. C **59**, 3009  
(1999) )

$\Lambda p \rightarrow \Lambda p$



- Production reactions with low momentum transfer



a)  $K^- d \rightarrow \pi^- p \Lambda$   
 (T.H. Tan Phys. Rev. Lett. **23**, 395 (1969))

$$a \approx a(S = 1) = -2.0 \pm 0.5 \text{ fm}$$

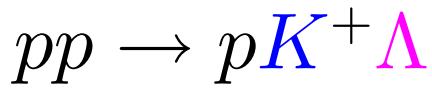
↑  
 FROM MODEL

b)  $K^- d \rightarrow \gamma n \Lambda$   
 ( R.L. Workman and W. Fearing Phys. Rev. C **41**, 1688 (1990); K.P. Gall, et al., Phys. Rev. C **42**, 475 (1990); W.R. Gibbs, Phys. Rev. C **61**, 064003 (2000))

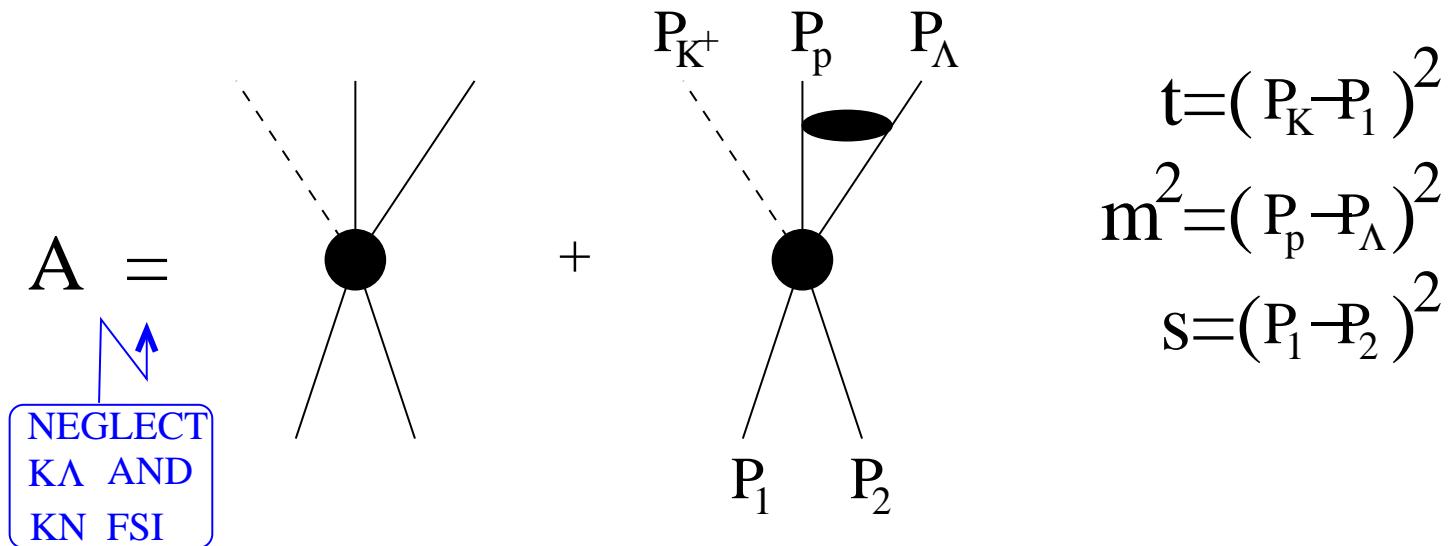
Problems:

- Theoretical errors not under control
- sensitive to short range  $\Lambda N$

- Production reactions  
with large momentum transfer



ASSUME POINT-LIKE PROD. OPERATOR



- IS A CONSTANT WITH RESPECT TO VARIATIONS IN  $m^2$

Dispersion relation technique :

$$A(s, t, m^2) = \exp \left[ \frac{1}{\pi} \int_{s_0}^{\infty} \frac{\delta(m'^2)}{m'^2 - m^2 - i0} dm'^2 \right] \times \Phi(s, t, m^2)$$

## Scattering length extraction

Effective range approximation:

$$\frac{1}{|D(k)|^2} = \frac{k^2 + \beta^2}{k^2 + \alpha^2}$$

$$a_S = \lim_{M^2 \rightarrow s_0} \frac{1}{2\pi} \left( \frac{m_\Lambda + m_p}{\sqrt{m_\Lambda m_p}} \right) \mathbf{P} \int_{s_0}^{s_{max}} dm^2 \sqrt{\frac{s_{max} - M^2}{s_{max} - m^2}} \\ \times \frac{1}{\sqrt{m^2 - s_0}(m^2 - M^2)} \log \left\{ \frac{1}{p'} \left( \frac{d^2 \sigma_S}{dm^2 dt} \right) \right\}$$

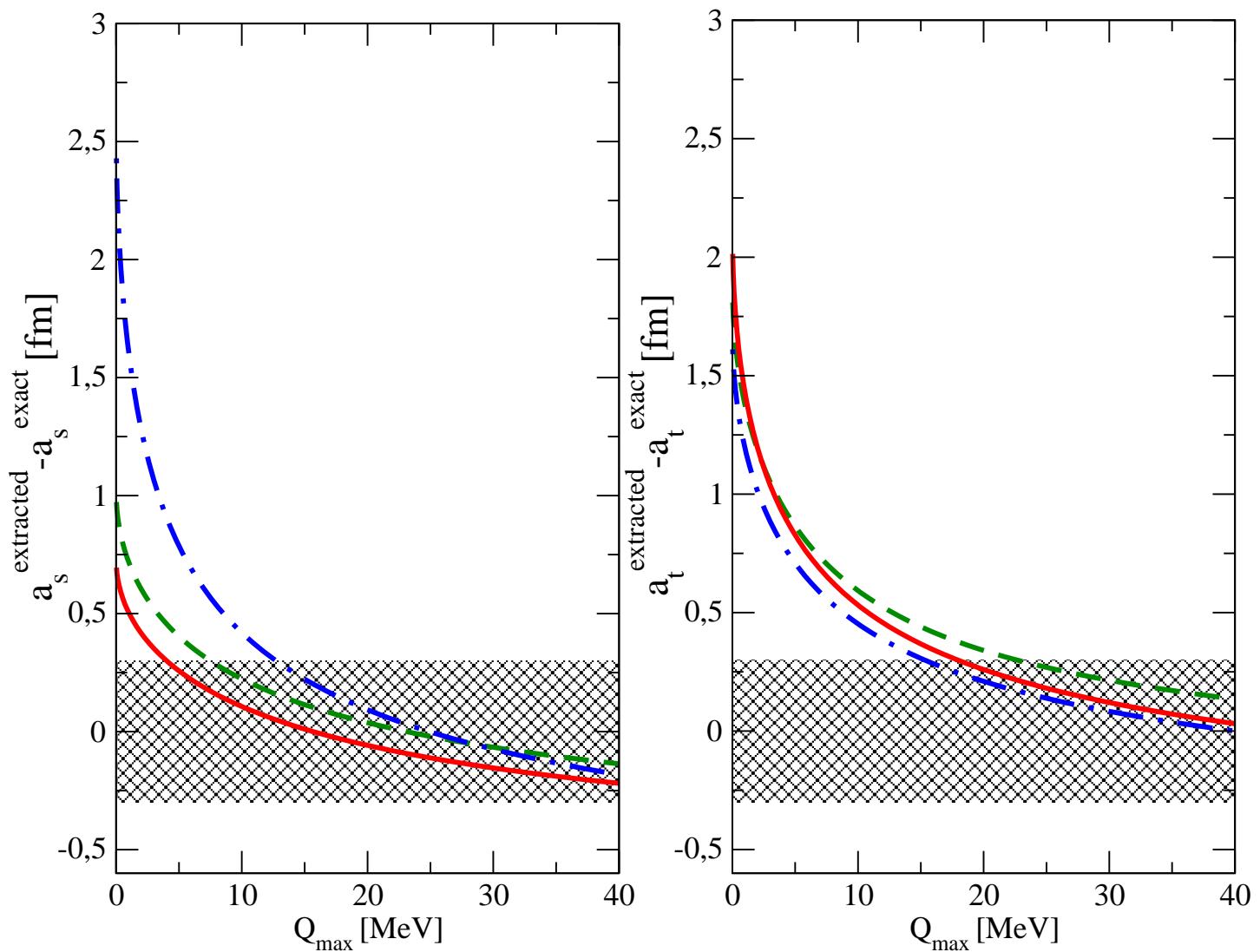
Relevant scales:  $a, r \Rightarrow Q_{max} \sim 20 - 40$  MeV

(N.I.Muskhelishvili, “*Singular Integral Equations*”, 1953;  
R. Omnes, Nuovo Cim. **8**, 316 (1958);  
W.R.Frazer and J.R.Fulco, Phys. Rev. Lett. **2**, 365 (1959);  
B.V.Geshkenbein, Yad.Fiz. **9**, 1232 (1969), Phys.Rev.D **61** )

# Comparison with model calculations

(A. Gasparian, Phys. Lett. B **480**, 273 (2000))

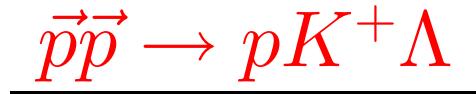
$$\Delta a \sim 0.3 \text{ fm}$$



— Nijmegen 99a

- - - Juelich 01

— · — Nijmegen 99f



Special kinematical conditions ( $q_{\Lambda N}$  is small)

+

parity conservation

+

Pauli principle (for initial protons)

⇓

$$A_{0y}\sigma_0(\theta = 90^\circ),$$

$$(1 - A_{xx})\sigma_0(\theta = 90^\circ \text{ or } \phi = 90^\circ) =$$

spin triplet contributions only,

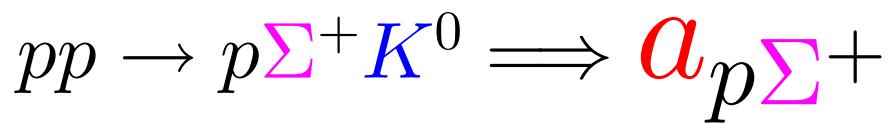
$$(1 + A_{xx} + A_{yy} - A_{zz})\sigma_0(\theta = 90^\circ) =$$

spin singlet contributions only.

## Further Applications

$D_{yy}\sigma_0(\theta = 90^\circ) =$   
triplet-singlet interference only.

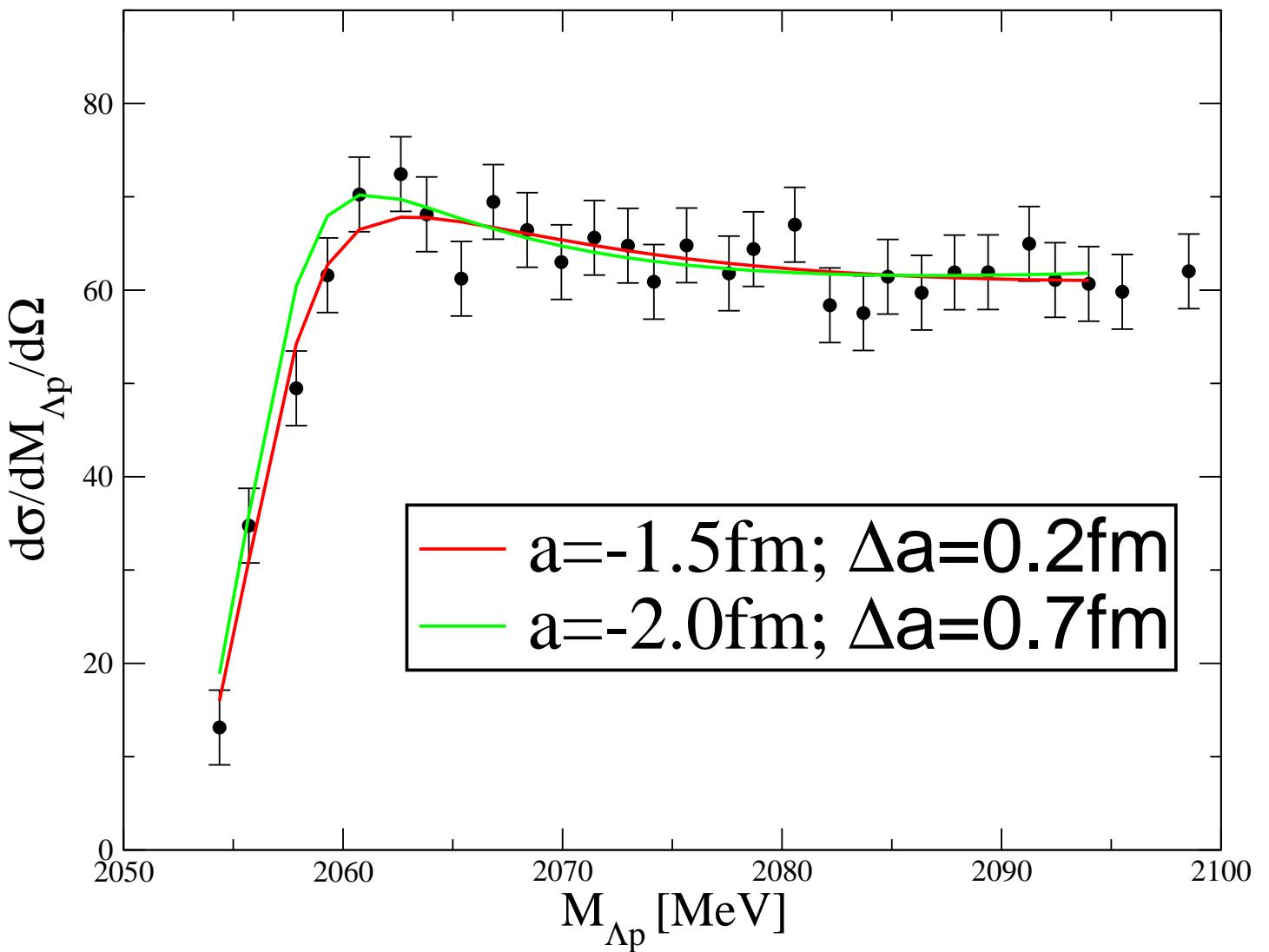
$$\frac{D_{yy}}{A_y} = Im \left\{ \exp \left[ \frac{1}{\pi} \int_{s_0}^{\infty} \frac{\delta_s(m'^2) - \delta_t(m'^2)}{m'^2 - m^2 - i0} dm'^2 \right] \times \Phi(s, t, m^2) \right\}$$



## Testing the method

$$\frac{d\sigma}{dM_{\Lambda p} d\Omega} = C_0 \exp \left\{ \frac{C_1}{C_2 - M_{\Lambda p}^2} + \dots \right\}$$

× Phase Space



## Summary

- A method for the determination of the  $\Lambda N$  scattering lengths in a model independent way is presented.
- Theoretical errors are under control.
- The polarization observables allow one to separate spin singlet and spin triplet states.
- $\frac{D_{yy}}{A_y}$  possibly allows one to extract  $a_s - a_t$ .
- $pp \rightarrow p\Sigma^+ K^0 \implies a_{p\Sigma^+}$ .
- A good mass resolution is needed to get a reasonable accuracy in the determination of the  $a_{\Lambda N}$ .